

## Lecture 04

# finite Square Well



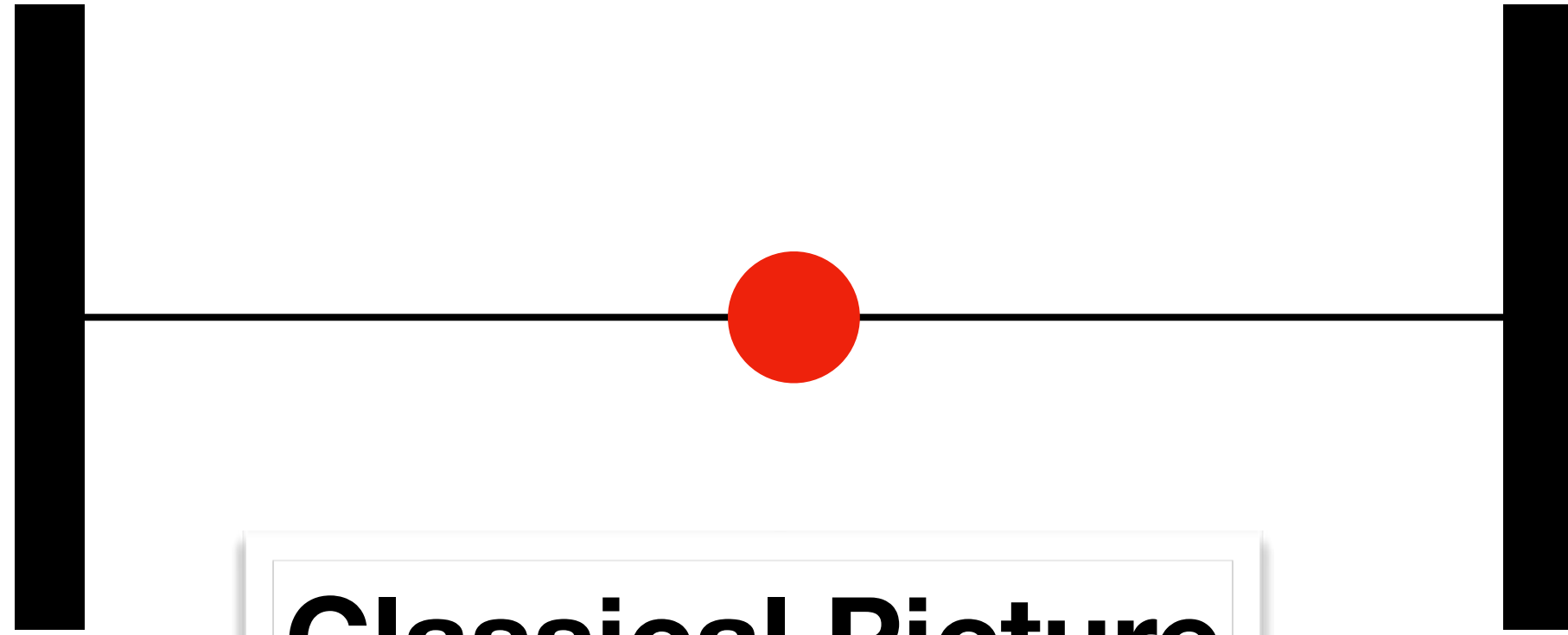
# Separation of the Time and Space Dependencies of $\psi(x, t)$

$$\frac{-\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$$

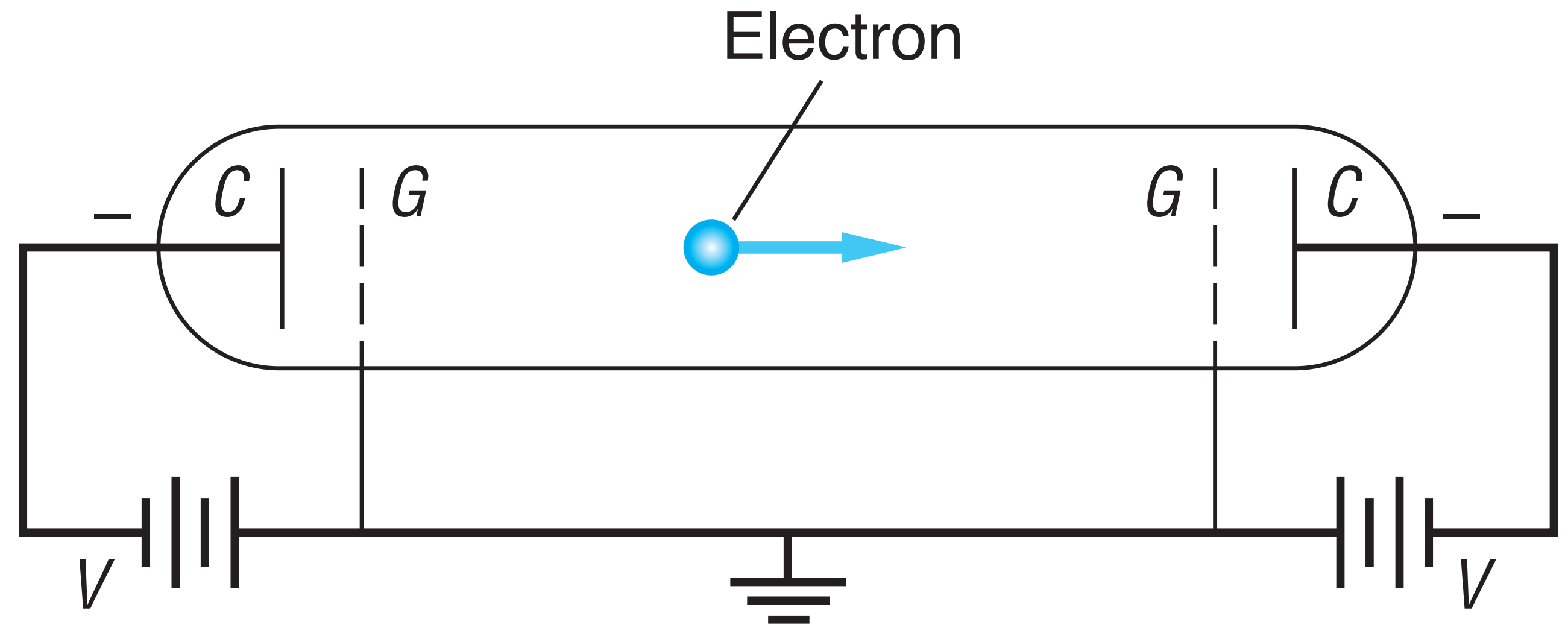
$$\Psi^*(x, t)\Psi(x, t) = \psi^*(x)e^{+iEt/\hbar}\psi(x)e^{-iEt/\hbar} = \psi^*(x)\psi(x)$$

1.  $\psi(x)$  must exist and satisfy the Schrödinger equation.
2.  $\psi(x)$  and  $d\psi/dx$  must be continuous.
3.  $\psi(x)$  and  $d\psi/dx$  must be finite.
4.  $\psi(x)$  and  $d\psi/dx$  must be single valued.
5.  $\psi(x) \rightarrow 0$  fast enough as  $x \rightarrow \pm\infty$  so that the normalization integral,

# Infinite potential well



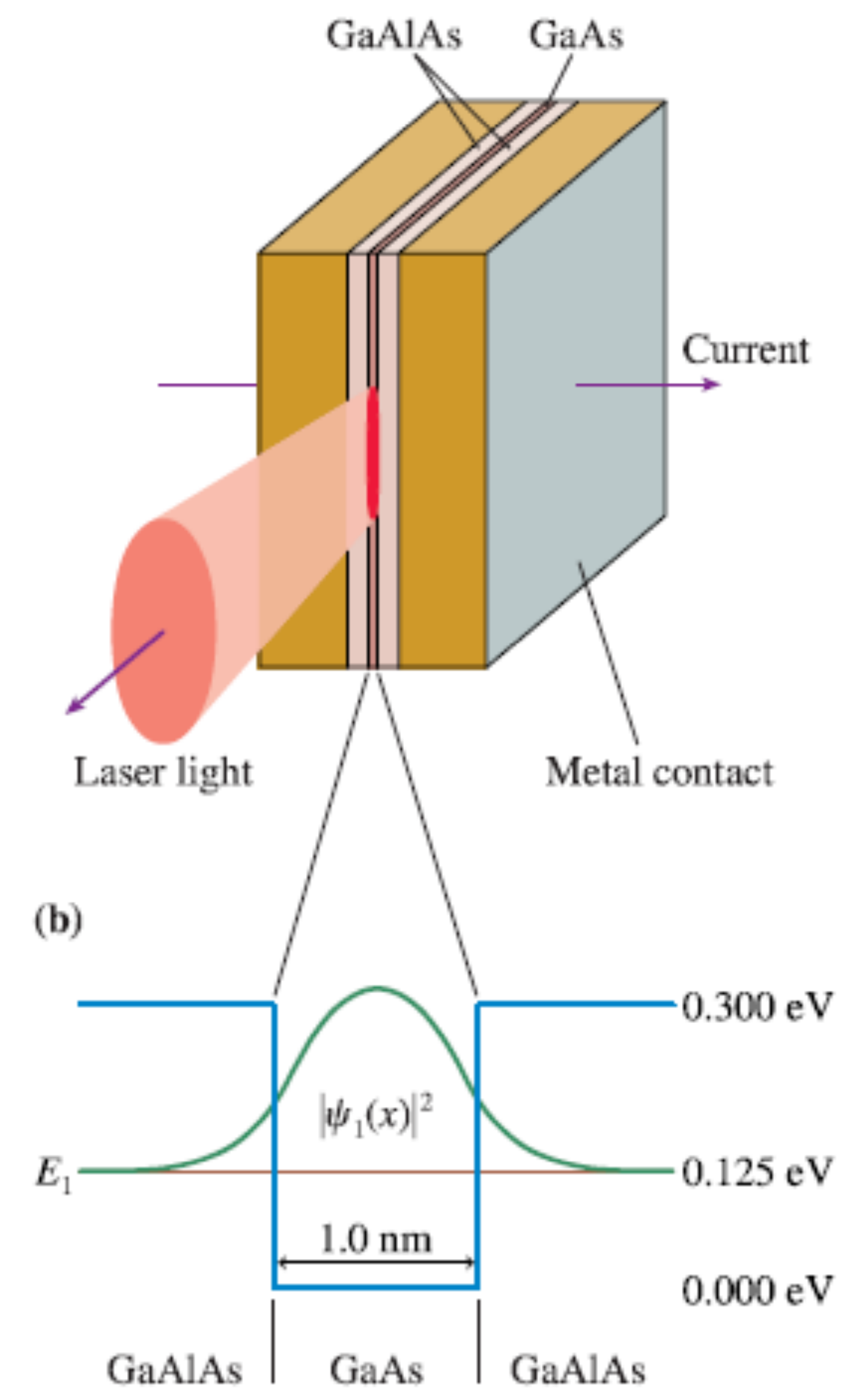
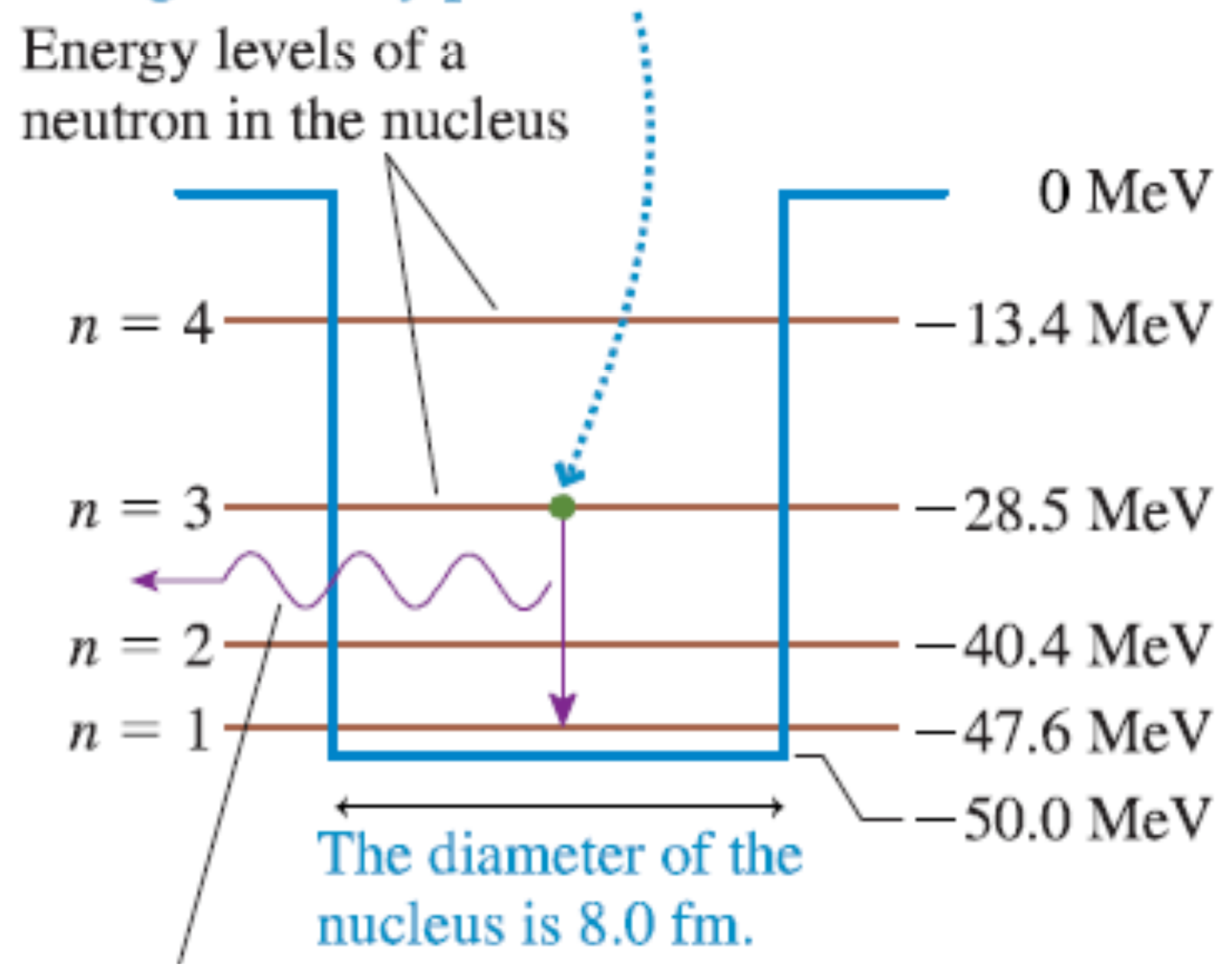
**Classical Picture**



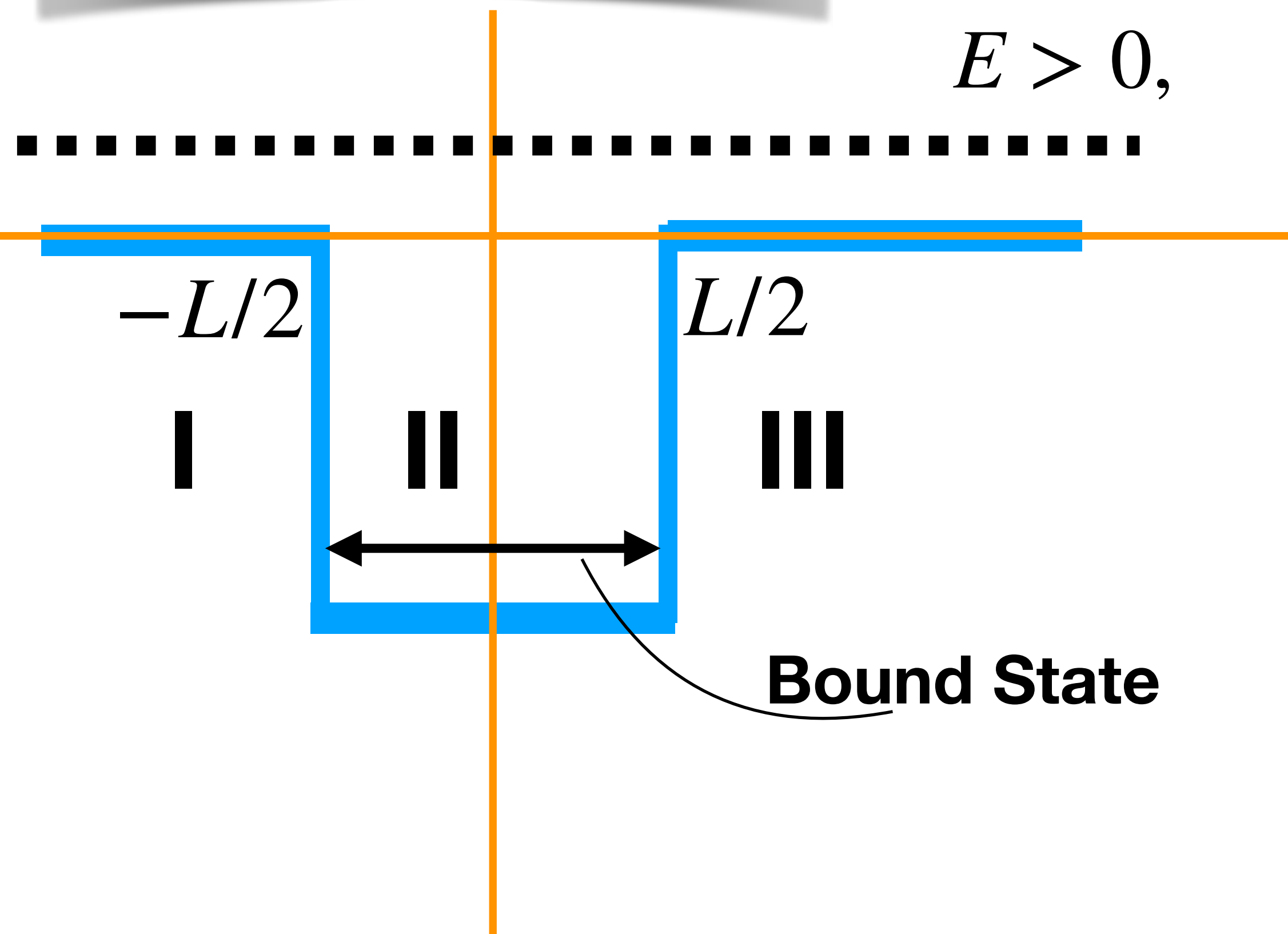
**Quantum Mechanical Realisation**



# Finite Potential Well



# Finite Potential Well



$$V(x) = 0, |x| > L$$

$$V(x) = -V_0, |x| < L$$

$$\Delta x \Delta p = \hbar$$

$$\Delta E = p^2/2m = \hbar^2/2mL^2$$

$$V_0/\Delta E \longrightarrow \text{Strength of Potential}$$



# Finite Potential Well



$$E > 0,$$

$$\Psi(x, t) = \phi(t)\psi(x)$$

$$\phi(t) = e^{-iE_n t/\hbar}$$

$$\frac{-\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$$

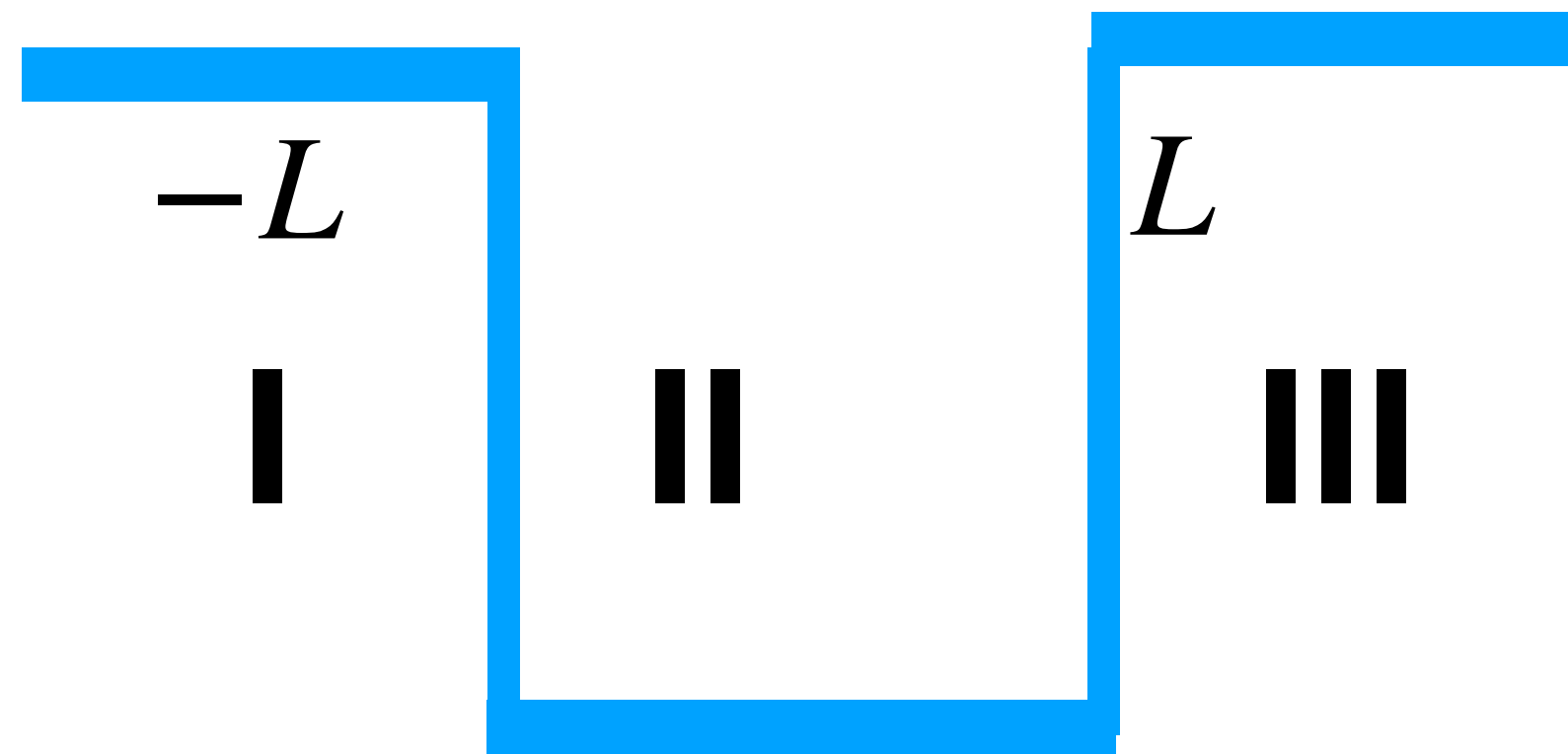
$$V(x) = 0 \quad \text{For region I and III}$$

$$\frac{-\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} = E\psi(x)$$

$$\frac{d^2\psi(x)}{dx^2} + \frac{2mE}{\hbar^2}\psi(x) = 0$$

$$k^2 = -\frac{2mE}{\hbar^2}$$

$$[k] = L^{-1}$$



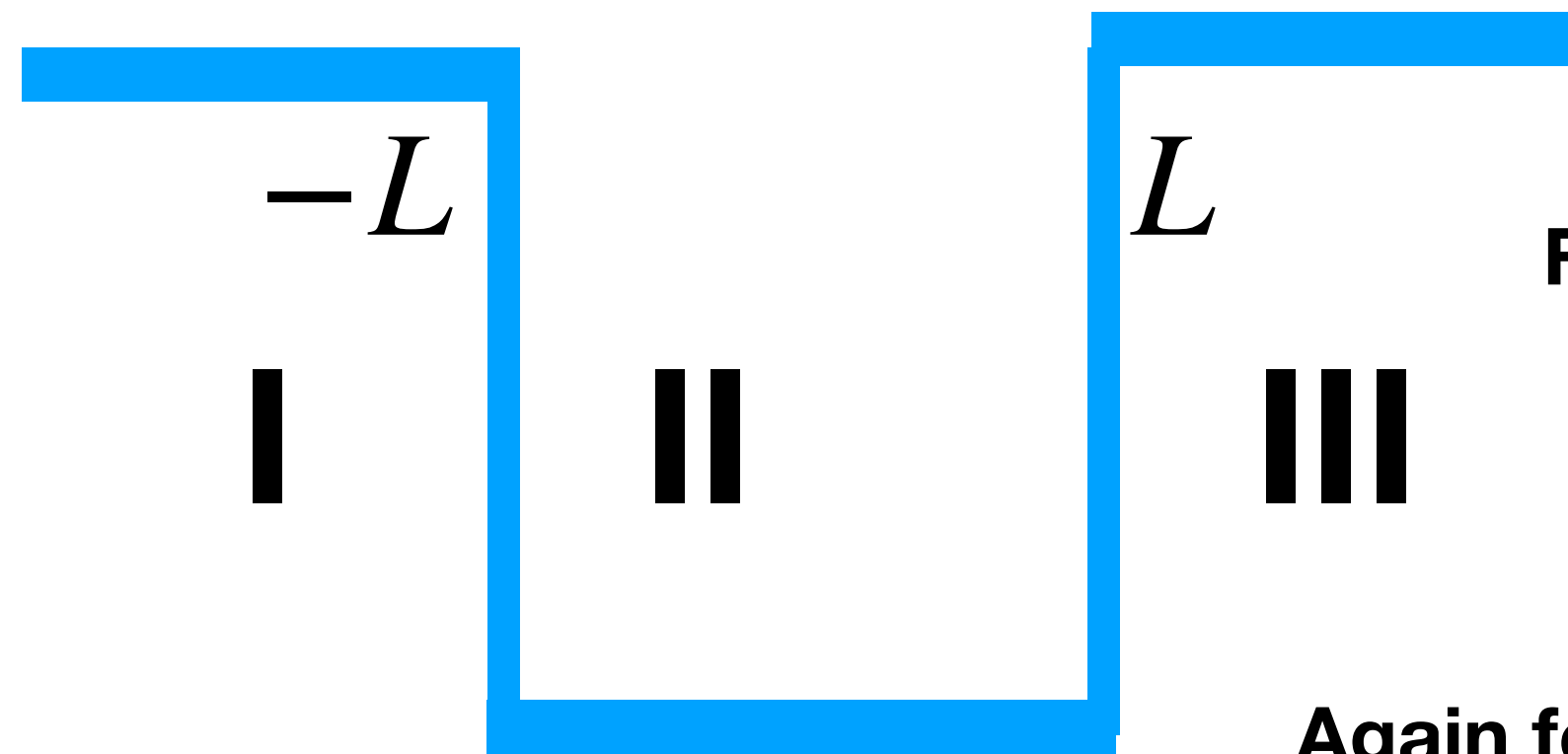


# Finite Potential Well



$$\psi_I(x) = Ae^{kx} + Be^{-kx}$$

$$\psi_{III}(x) = Ee^{kx} + Fe^{-kx}$$



For wave functions to make sense  $B=0$  and  $E=0$

$$\psi_I(x) = Ae^{kx}$$

$$\psi_{III}(x) = Fe^{-kx}$$

Again for region II

$$\begin{aligned} \frac{-\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} - V_0\psi(x) &= E\psi(x) \\ \frac{d^2\psi(x)}{dx^2} &= -\frac{2m}{\hbar^2} [V_0 + E]\psi(x) \end{aligned}$$

Now Let  $k'^2 = \frac{2m}{\hbar^2} (V_0 + E)$

$$\frac{d^2\psi(x)}{dx^2} = -k'^2\psi(x)$$

$$\psi_{II}(x) = C\sin k'x + D\cos k'x$$

# Finite Potential Well



$$\psi_{II}(x) = C\sin k'x + D\cos k'x$$

$$\psi_I(x) = Ae^{kx}$$

$$\psi_{III}(x) = Fe^{-kx}$$

1.  $\psi(x)$  must exist and satisfy the Schrödinger equation.
2.  $\psi(x)$  and  $d\psi/dx$  must be continuous.
3.  $\psi(x)$  and  $d\psi/dx$  must be finite.
4.  $\psi(x)$  and  $d\psi/dx$  must be single valued.
5.  $\psi(x) \rightarrow 0$  fast enough as  $x \rightarrow \pm\infty$  so that the normalization integral,

$$\psi_I(-L) = \psi_{II}(-L) \quad 1$$

$$\frac{d\psi_I(-L)}{dx} = \frac{d\psi_{II}(-L)}{dx} \quad 2$$

$$\psi_{II}(L) = \psi_{III}(L) \quad 3$$

$$\frac{d\psi_{II}(L)}{dx} = \frac{d\psi_{III}(L)}{dx} \quad 4$$



# Finite Potential Well

$$Ae^{-Lk} = -C\sin k'L + D\cos k'L \quad 5$$

$$-Ake^{-Lk} = -Ck'\cos k'L - Dk'\sin k'L \quad 6$$

$$Fe^{Lk} = C\sin k' + D\cos k'L \quad 7$$

$$Fke^{Lk} = Ck'\cos k'L - Dk'\sin k'L \quad 8$$

**Even parity**  $\psi(-x) = \psi(x)$

$$Fe^{-Lk} = D\cos k'L$$

$$-Fke^{-Lk} = -LDk\sin k'L$$

$$k = k'\tan k'L$$



# Finite Potential Well



$$k^2 = -\frac{2mE}{\hbar^2} \quad k'^2 = \frac{2m}{\hbar^2}(V_0 + E)$$

$$k^2 + k'^2 = \frac{2mV_0}{\hbar^2}$$

$$k'L = \zeta \quad \zeta_0 = \frac{L}{\hbar} \sqrt{2mV_0}$$

$$\zeta^2 + (kL)^2 = \zeta_0^2$$

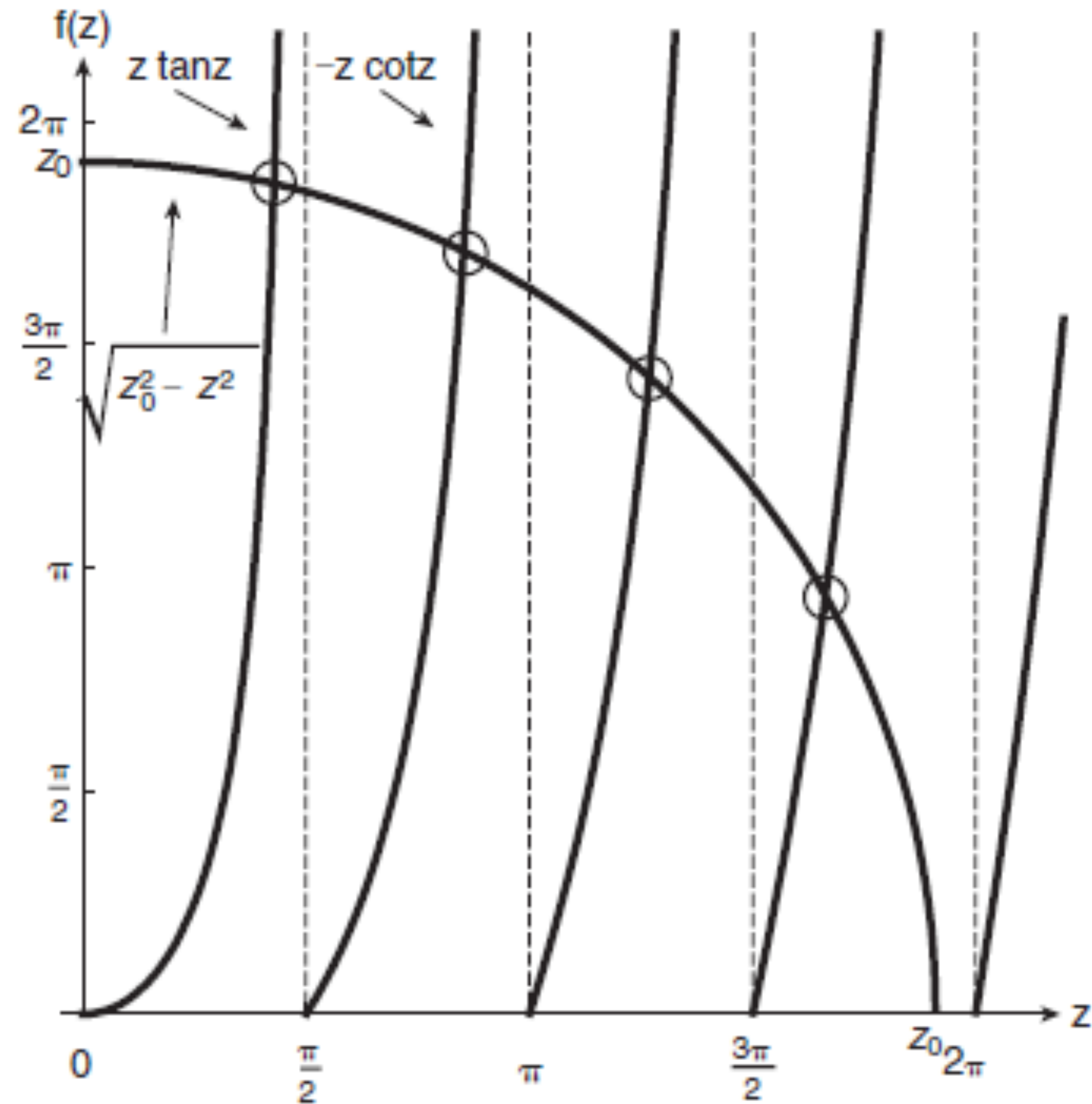
$$kL = \sqrt{\zeta_0^2 - \zeta^2}$$

$$k = k' \tan k'L$$

$$kL = k' L \tan k'L$$

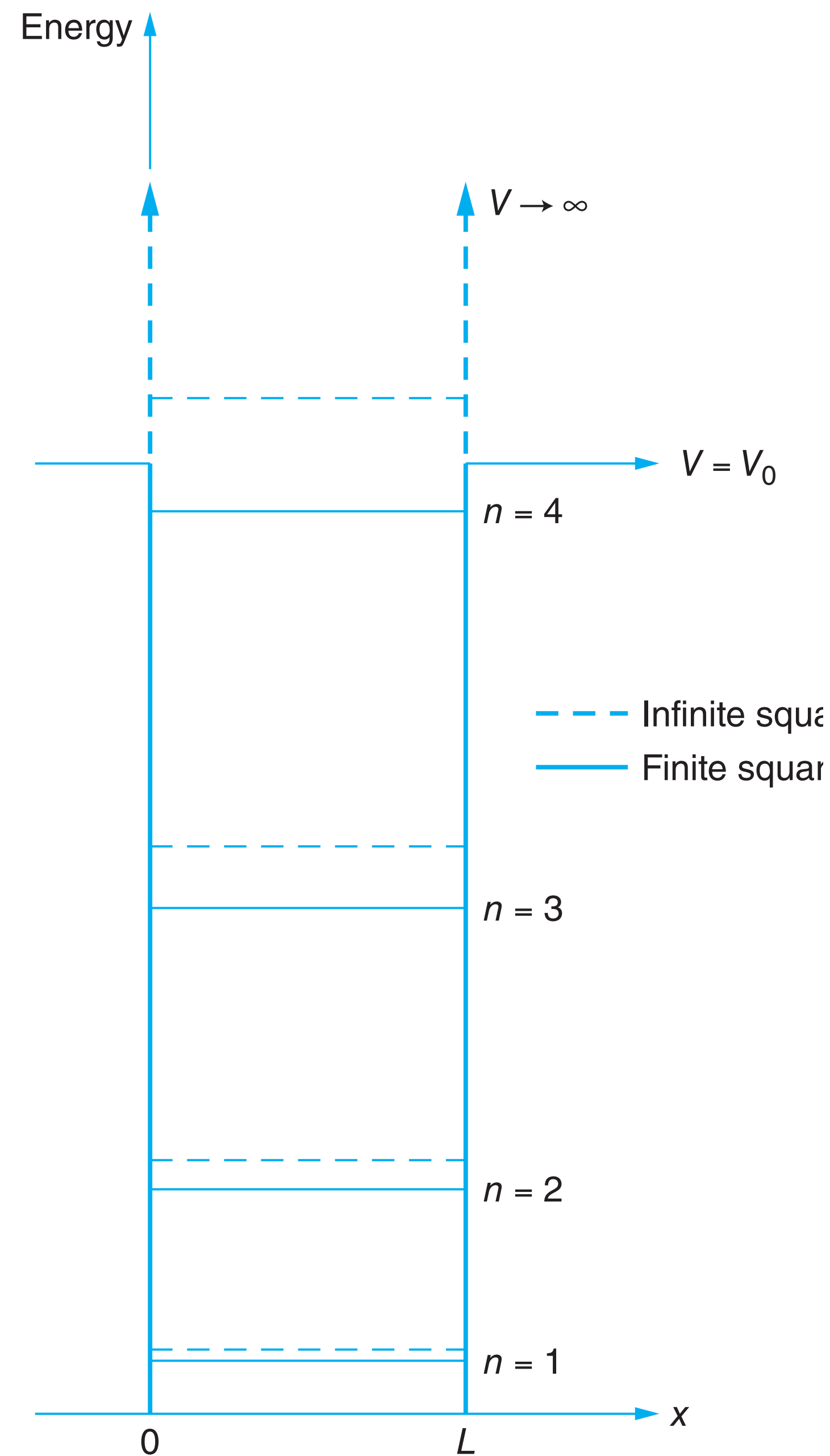
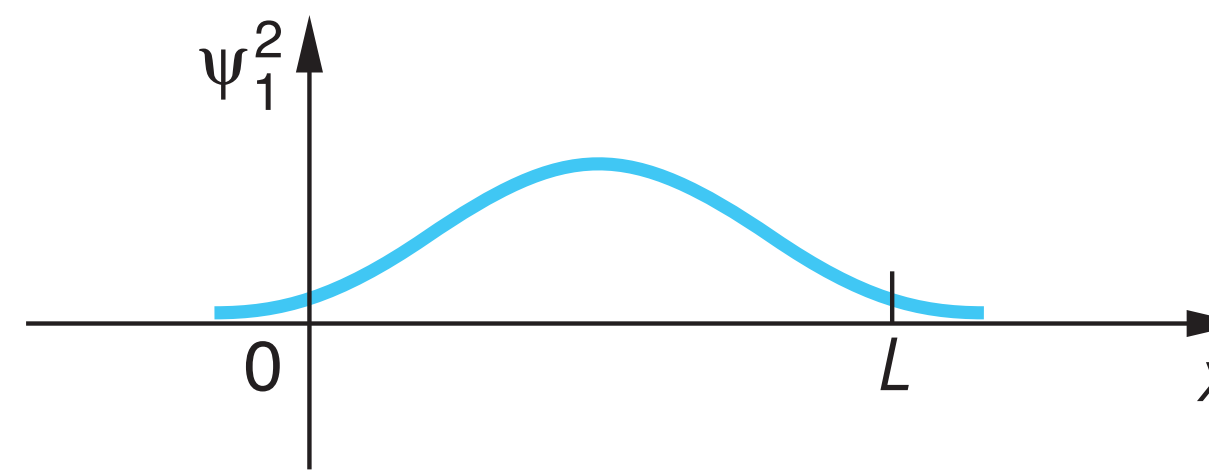
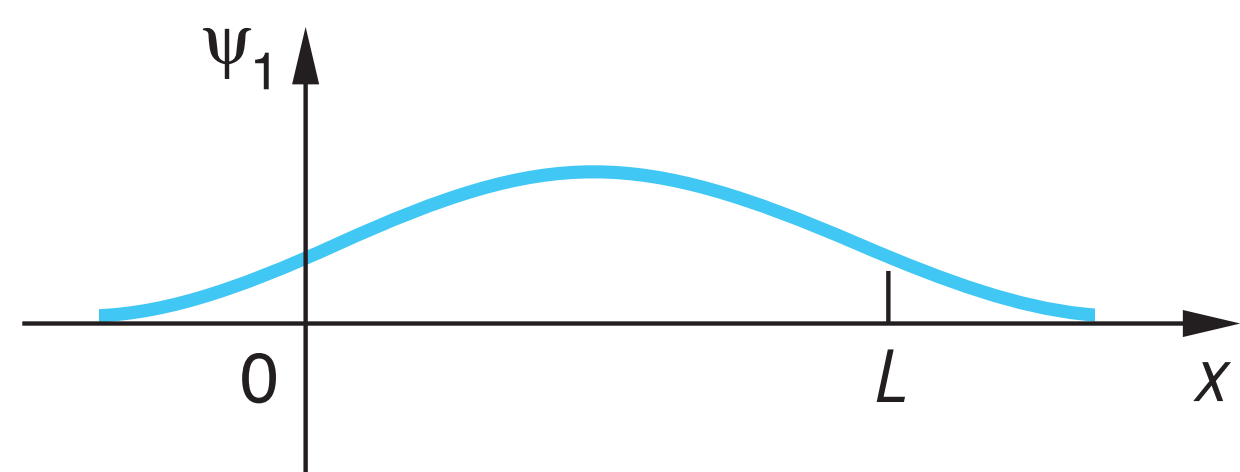
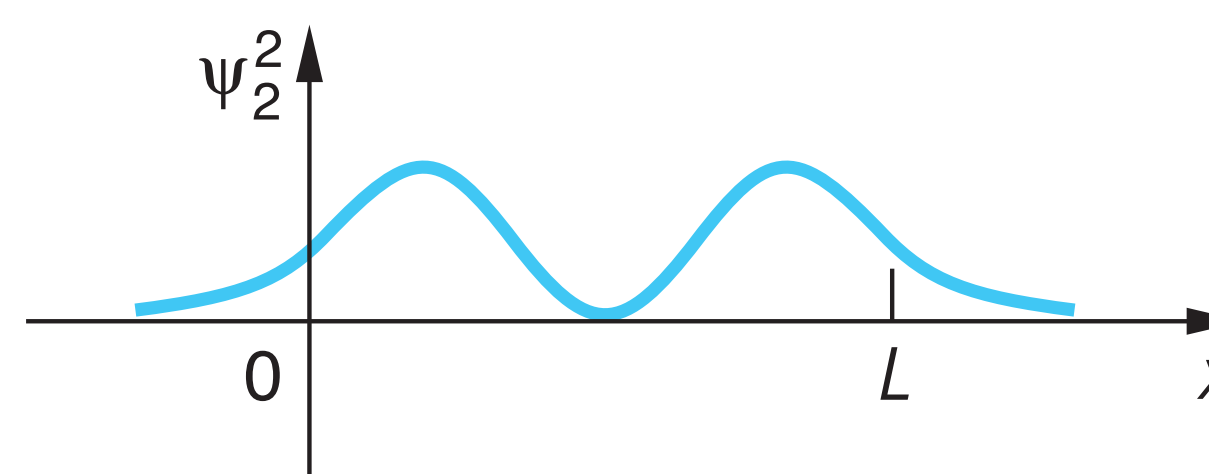
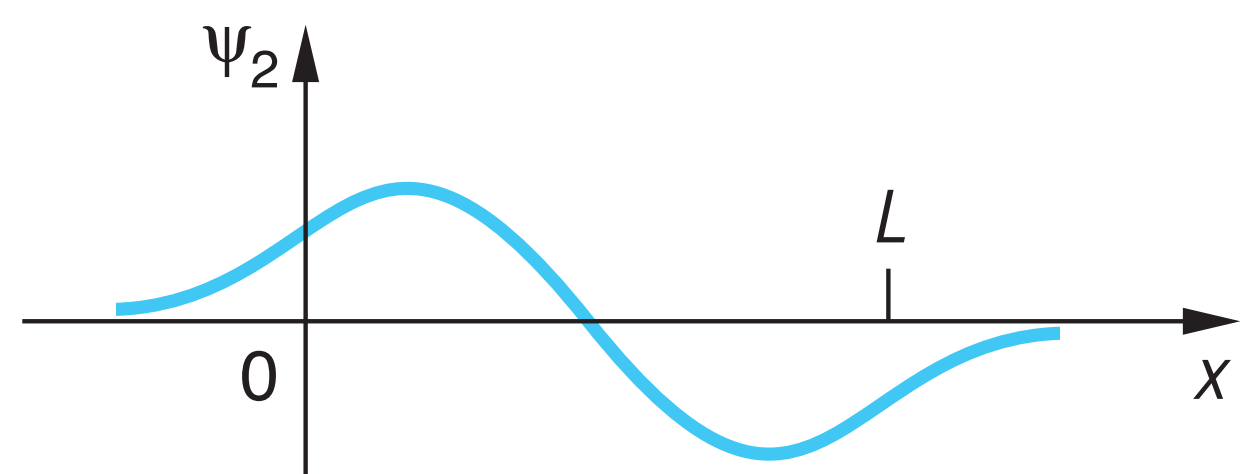
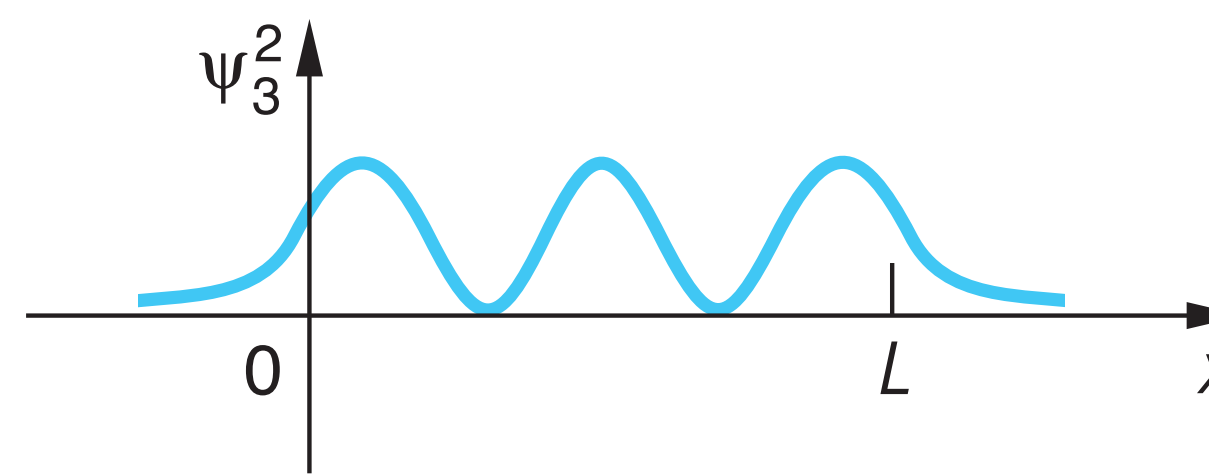
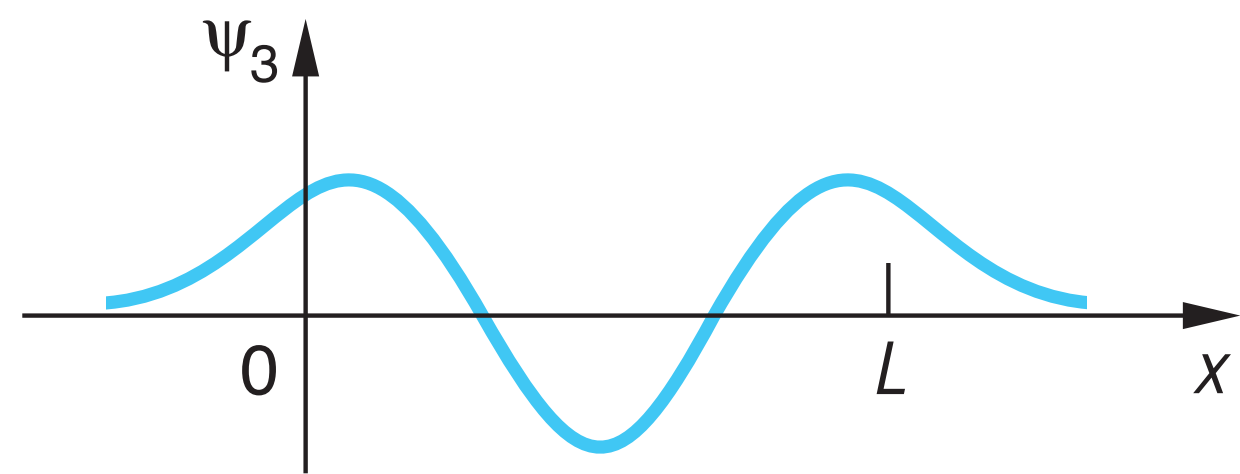
$$\zeta \tan \zeta = \sqrt{\zeta_0^2 - \zeta^2}$$

# Finite Potential Well

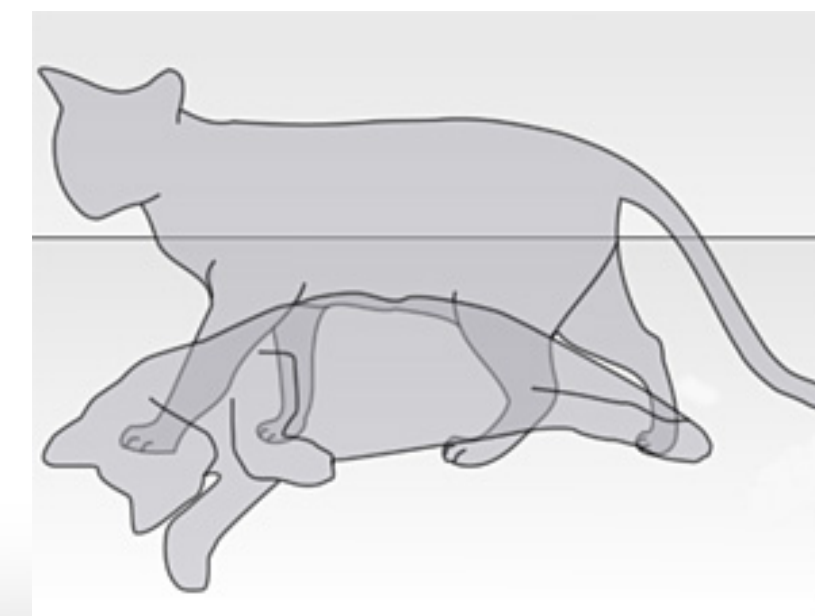


Graphical solution of the transcendental equations for the allowed energies of a finite square well ( $z_0 = 6$ ).

# Finite Potential Well







# Curiosity Kills the Cat

*Lecture 04*  
*Concluded*